

Implementation of FEM on HPC – I

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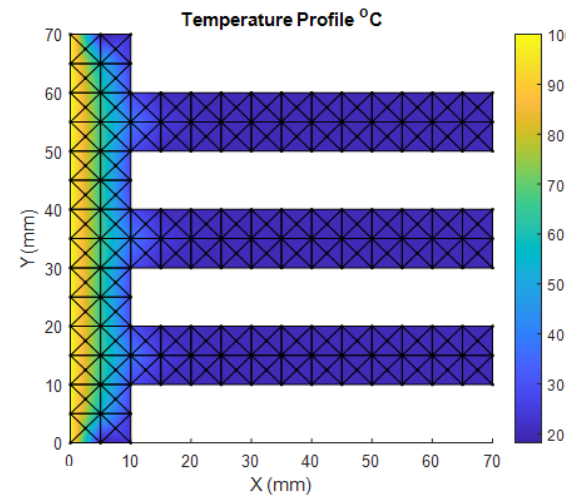
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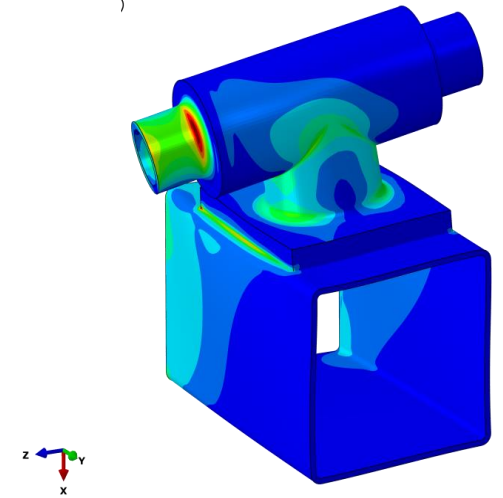
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Methods for solution of Boundary Value Problem (BVP):

- Finite Difference Method (FDM)
- Finite Element Method (FEM)
- Boundary Element Method (BEM)
- Finite Volume Method (FVM)



$$k \Delta T(x, y) + q_v = 0$$



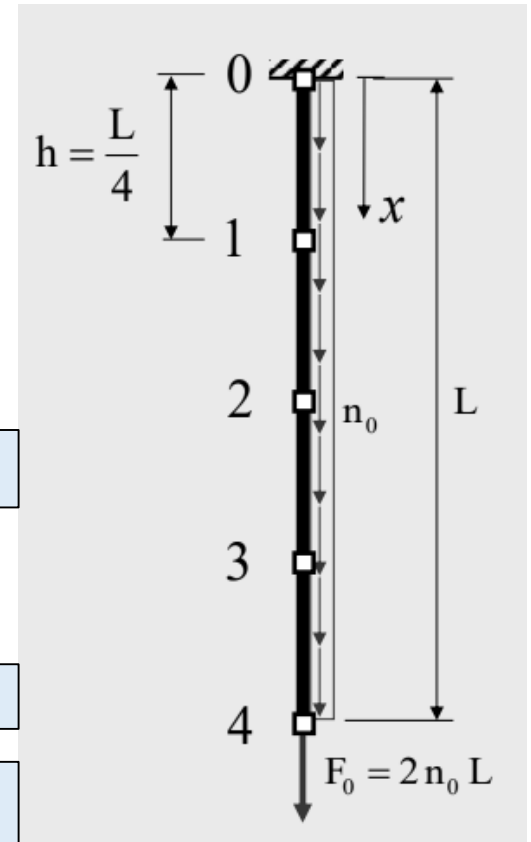
$$\begin{aligned} \vec{\nabla}^T \sigma + \mathbf{b} &= \mathbf{0} \\ \sigma &= \mathbf{D}\epsilon - \mathbf{D}\epsilon_0 \\ \epsilon &= \vec{\nabla} \mathbf{u} \end{aligned}$$

- IDEA: approximation of derivatives in governing equation and boundary conditions at points (grid)
- Approximation is derived based on the Taylor's series expansion
- ADVANTAGE: simple to use
- DISADVANTAGE: irregular domains, mesh preparation

$$EA u''(x) = -n_0$$

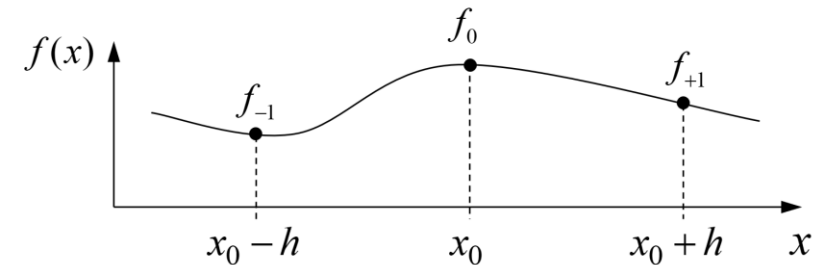
$$u(0) = 0$$

$$EA u(L)' = F_0$$



- Approximation based on the Taylor's series expansion

$$f(x_0 + h) = f(x_0) + \frac{h}{1!} \frac{d^1 f(x_0)}{dx^1} + \frac{h^2}{2!} \frac{d^2 f(x_0)}{dx^2} + \frac{h^3}{3!} \frac{d^3 f(x_0)}{dx^3} + \dots$$



$$f_{+1} \approx f_0 + h f'_0 + \frac{h^2}{2} f''_0$$

$$f'_0 \approx \frac{f_{+1} - f_{-1}}{2h}$$

$$f_{-1} \approx f_0 - h f'_0 + \frac{h^2}{2} f''_0$$

$$f''_0 \approx \frac{f_{+1} - 2f_0 + f_{-1}}{h^2}$$

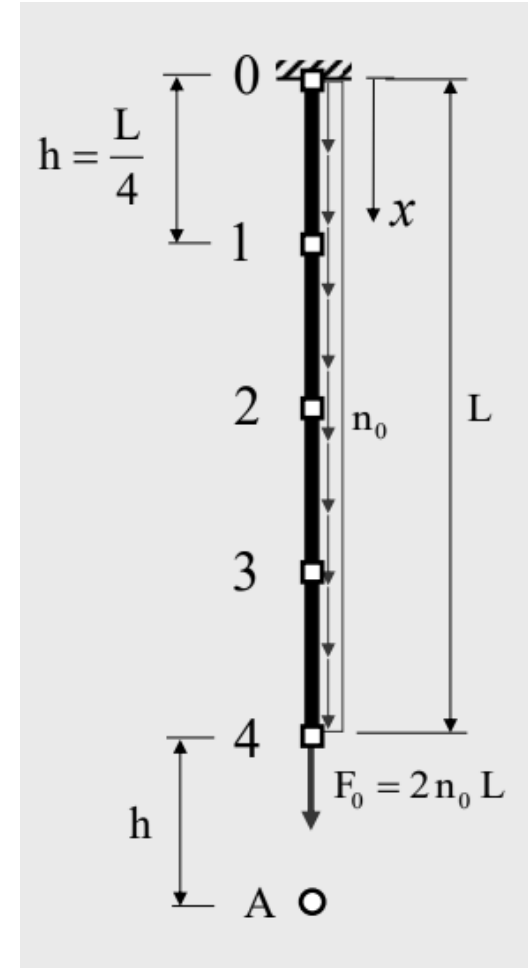
Finite difference approximation
of 1st and 2nd derivative

- Writing boundary conditions in discrete form:

$$u(0)=0 \Rightarrow u_0 = 0 \quad EA u(L)' = F_0 \Rightarrow u_A - u_3/(2h) = F_0$$

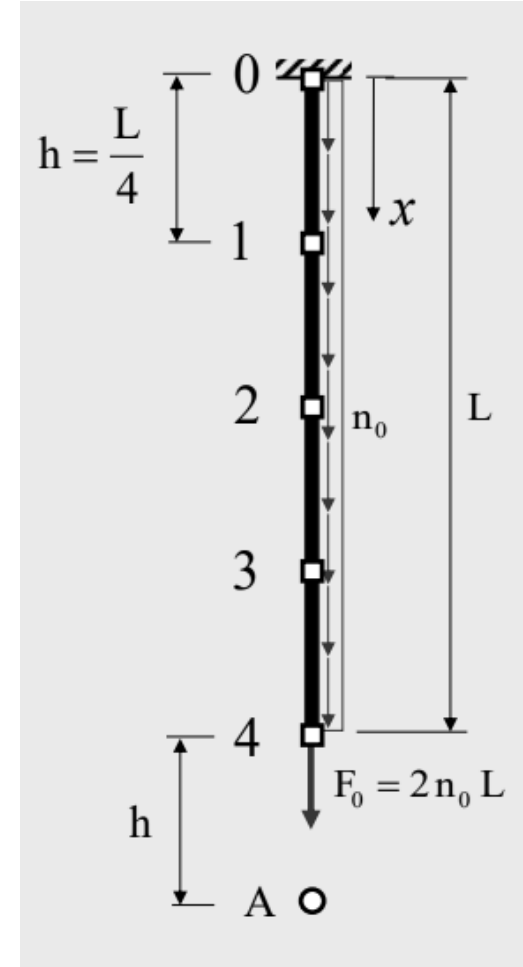
- Writing governing equation at discrete points 1 to 4

$$EA u''(x) = -n_0 \Rightarrow \begin{aligned} u_0 - 2u_1 + u_2 &= -n_0 h^2 / (EA) \\ u_1 - 2u_2 + u_3 &= -n_0 h^2 / (EA) \\ u_2 - 2u_3 + u_4 &= -n_0 h^2 / (EA) \\ u_3 - 2u_4 + u_A &= -n_0 h^2 / (EA) \end{aligned}$$

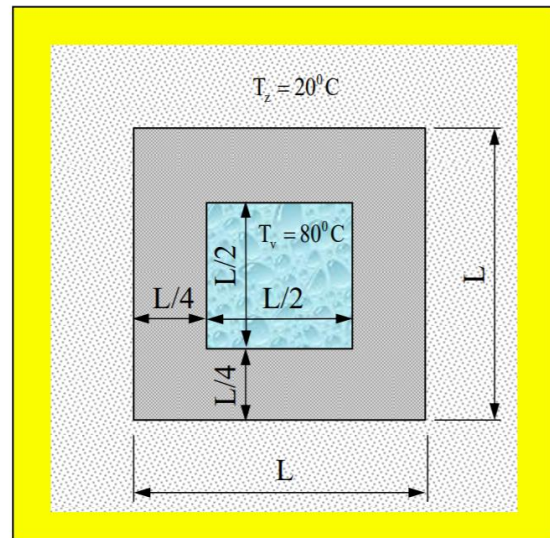


- System of equations:

	u_0	u_1	u_2	u_3	u_4	u_A	
BC pt. 0 →	1	0	0	0	0	0	$\left. \begin{matrix} u_0 \\ u_1 \\ u_2 \\ u_3 \\ u_4 \\ u_A \end{matrix} \right\} = \left\{ \begin{matrix} 0 \\ 16 \\ -1 \\ -1 \\ -1 \\ -1 \end{matrix} \right\} \frac{n_0 L^2}{16EA}$
BC pt. 4 →	0	0	0	-1	0	+1	
GE pt. 1 →	+1	-2	+1	0	0	0	
GE pt. 2 →	0	+1	-2	+1	0	0	
GE pt. 3 →	0	0	+1	-2	+1	0	
GE pt. 4 →	0	0	0	+1	-2	+1	

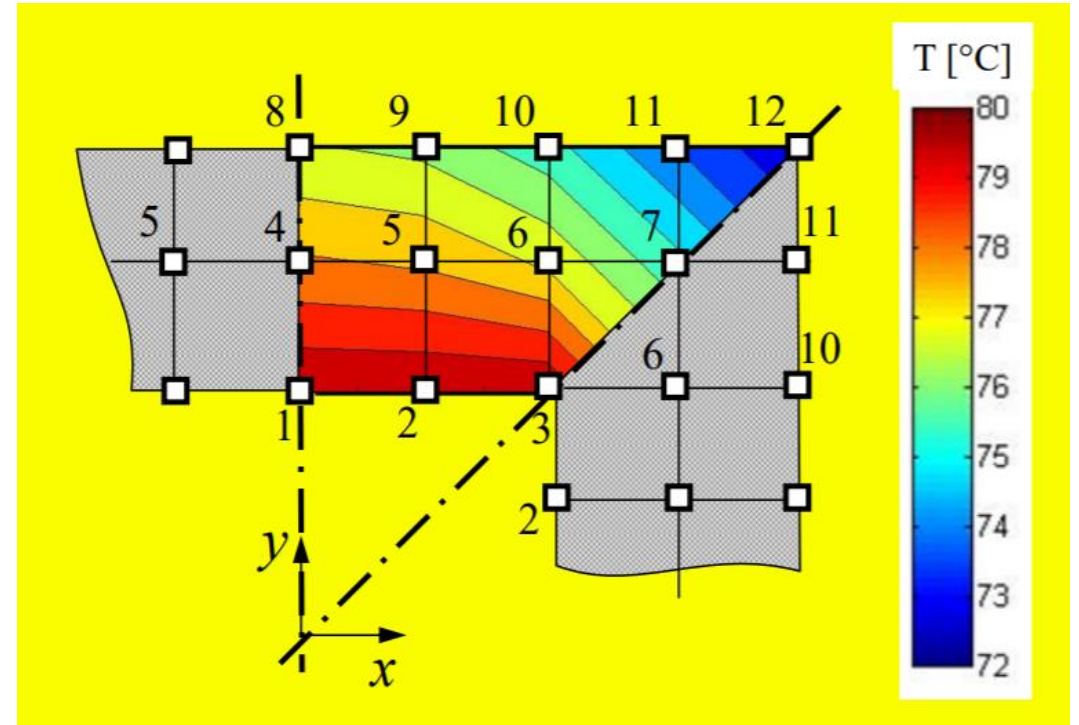
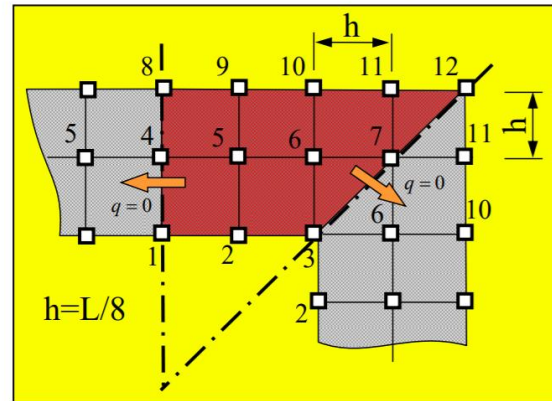
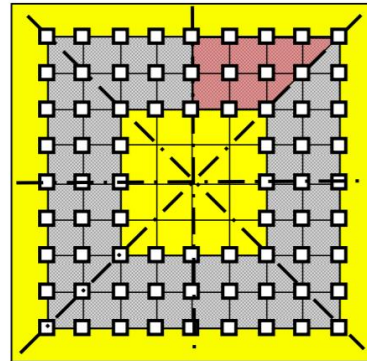


- Stationary heat transfer in 2D



MNM: XIV/19/3f

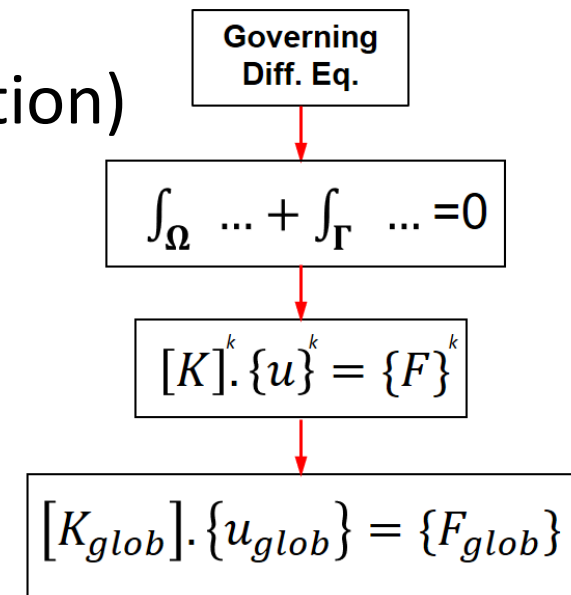
$$k \Delta T(x, y) + q_v = 0$$



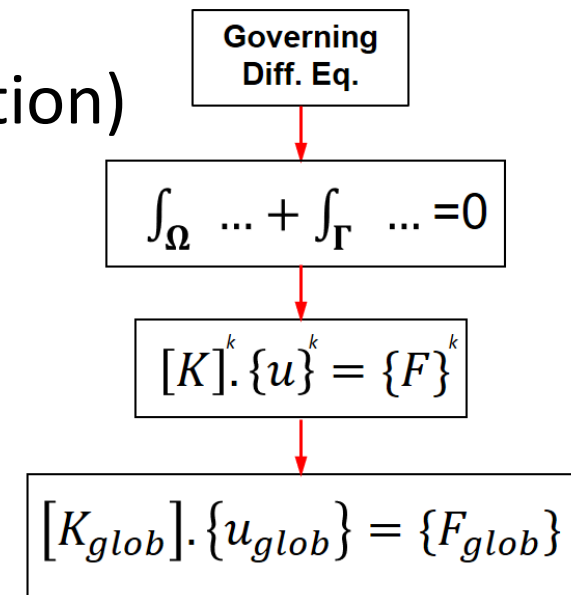
- IDEA: integral based approach
- Starting point of derivation is Weak Integral Form of governing equation
- Solution domain is discretized into sub-domains, called finite elements (FEs)
- In the FE sub-domain we approximate unknown quantities

- ADVANTAGE: simple to use for complex geometrical domains, useful for all types of physical problems
- DISADVANTAGE: computationally intensive method

- Derivation of Weak Integral Form of Governing Equation
- Approximation of Solution Variable over Finite Element Domain
- Derivation of Finite Element Matrix Equation
- Meshing & Writing FE Matrix Equation for each FE
- The Assembly Process (Global Finite Element Matrix Equation)
- Writing Boundary Conditions
- Solution of System of Equations



- **Derivation of Weak Integral Form of Governing Equation**
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Derivation of Weak Integral Form of Governing Equation

$$EA u''(x) = -n_0 \quad \text{Governing Equation}$$

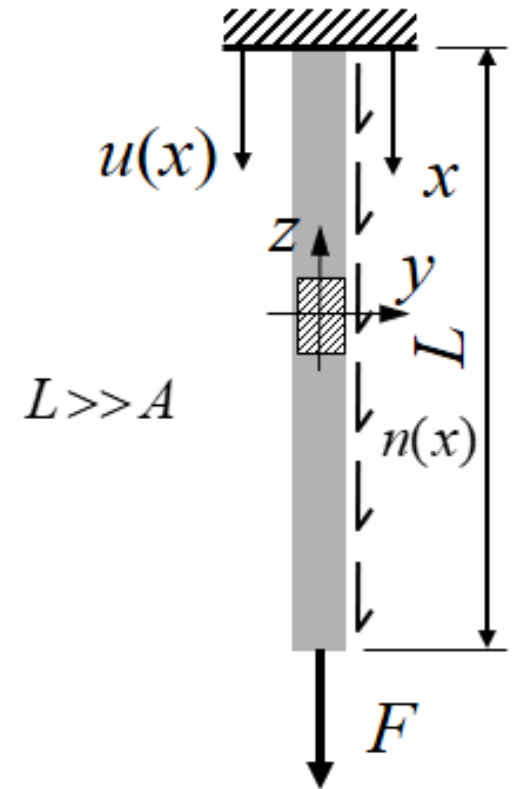


$$EA u''(x) + n_0 = 0$$



$$\int_0^{L_e} (EA u''(x) + n_0)v(x)dx = 0$$

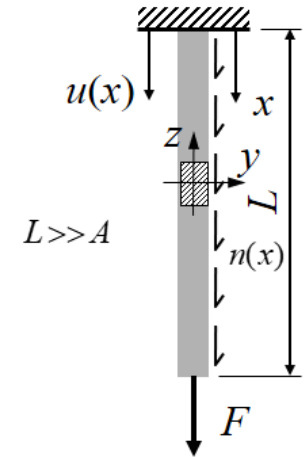
Strong Form



Derivation of Weak Integral Form of Governing Equation

$$\int_0^{L_e} EA u''(x) v(x) dx = - \int_0^{L_e} n_0 v(x) dx$$

Strong Form



Integration by PARTS 

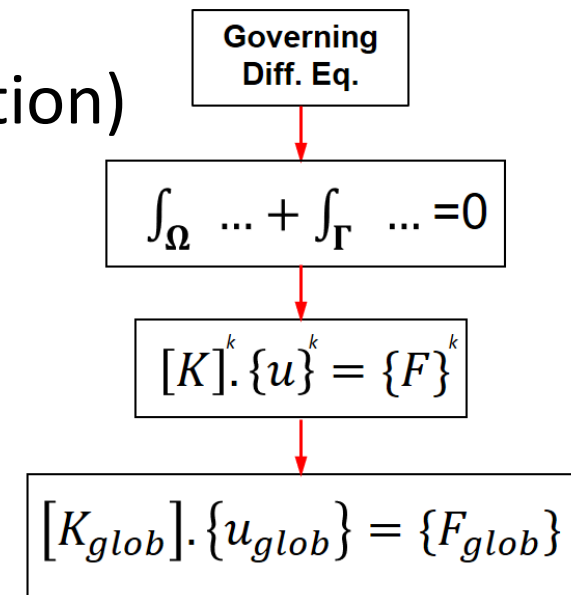
$$\int_0^{L_e} EA u'(x) v'(x) dx = EA u'(L) v(L) - EA u'(0) v(0) + \int_0^{L_e} n_0 v(x) dx$$

WEAK Form


 $N(L)$


 $N(0)$

- Derivation of Weak Integral Form of Governing Equation
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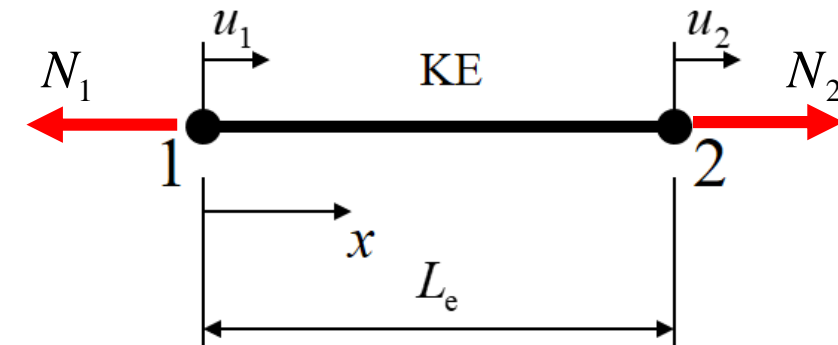
$$\int_0^{L_e} EA u'(x)v'(x)dx = N_2 v(L) - N_1 v(0) + \int_0^{L_e} n_0 v(x)dx$$

Approximation of $u(x)$

$$u(x) \approx \hat{u}(x) = u_1 \psi_1(x) + u_2 \psi_2(x)$$

$$\hat{u}(0) = u_1 \Rightarrow \psi_1(0) = 1 \wedge \psi_1(L_e) = 0 \Rightarrow \psi_1(x) = 1 - \frac{x}{L_e}$$

$$\hat{u}(L_e) = u_2 \Rightarrow \psi_2(0) = 0 \wedge \psi_2(L_e) = 1 \Rightarrow \psi_2(x) = \frac{x}{L_e}$$



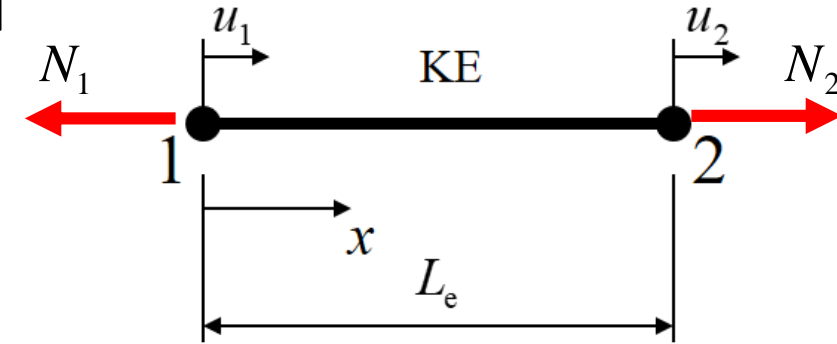
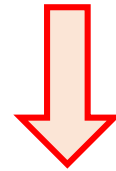
Shape functions

FEM – Matrix Equation

$$\int_0^{L_e} EA u'(x) v'(x) dx = N_2 v(L) - N_1 v(0) + \int_0^{L_e} n_0 v(x) dx$$

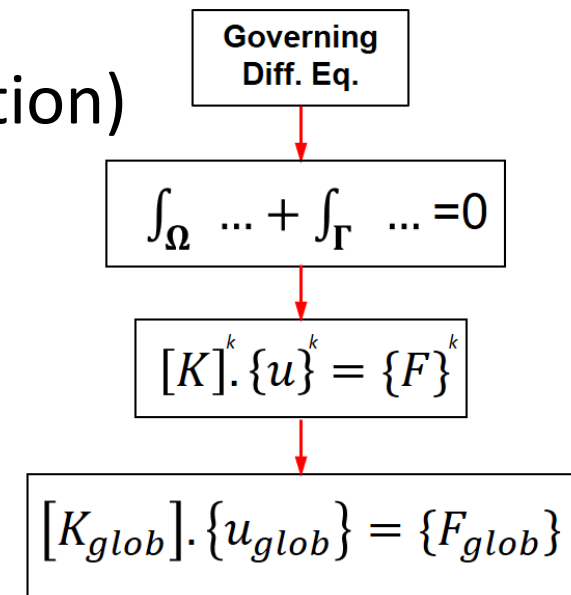
Inserting $\hat{u}(x)$ into weak form

$$u'(x) \approx \hat{u}'(x) = u_1 \frac{-1}{L_e} + u_2 \frac{1}{L_e}$$



$$EA \int_0^{L_e} u'(x) v'(x) dx = EA \int_0^{L_e} \left[u_1 \frac{-1}{L_e} + u_2 \frac{1}{L_e} \right] v'(x) dx = N_2 v(L) - N_1 v(0) + \int_0^{L_e} n_0 v(x) dx$$

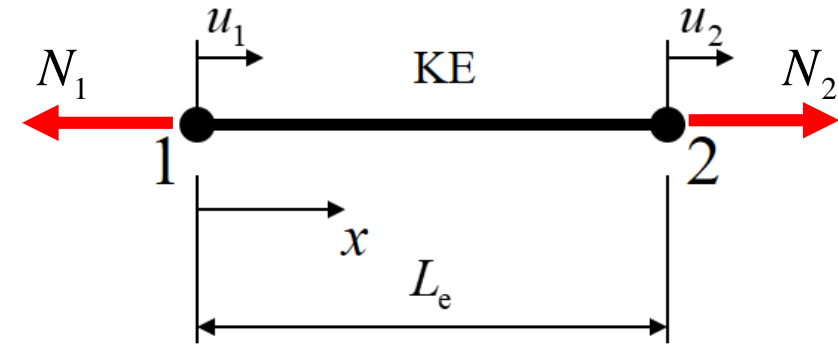
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$$EA \int_0^{L_e} u'(x) v'(x) dx = EA \int_0^{L_e} \left[\left(u_1 \frac{-1}{L_e} + u_2 \frac{1}{L_e} \right) \right] v'(x) dx = N_2 v(L) - N_1 v(0) + \int_0^{L_e} n_0 v(x) dx$$

Choice of $v(x)$ – Galerkin approach

① $v(x) = \psi_1(x) = 1 - \frac{x}{L_e}, \quad v'(x) = -\frac{1}{L_e}$



$$EA \int_0^{L_e} \left[\left(u_1 \frac{-1}{L_e} + u_2 \frac{1}{L_e} \right) \right] \left(\frac{-1}{L_e} \right) dx = -N_1 + \int_0^{L_e} n_0 \psi_1(x) dx$$



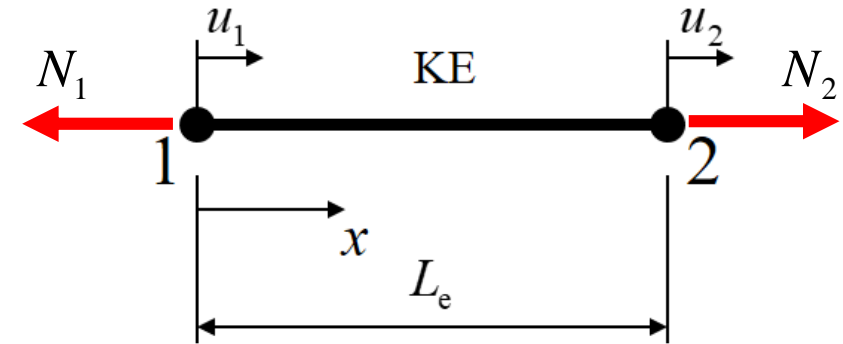
$$\frac{EA}{L_e} (u_1 - u_2) = -N_1 + \int_0^{L_e} n_0 \psi_1(x) dx$$

FEM – Matrix Equation

$$EA \int_0^{L_e} u'(x) v'(x) dx = EA \int_0^{L_e} \left[\left(u_1 \frac{-1}{L_e} + u_2 \frac{1}{L_e} \right) \right] v'(x) dx = N_2 v(L) - N_1 v(0) + \int_0^{L_e} n_0 v(x) dx$$

Choice of $v(x)$ – Galerkin approach

② $v(x) = \psi_2(x) = \frac{x}{L_e}, \quad v'(x) = \frac{1}{L_e}$



$$EA \int_0^{L_e} \left[\left(u_1 \frac{-1}{L_e} + u_2 \frac{1}{L_e} \right) \right] \left(\frac{1}{L_e} \right) dx = N_2 + \int_0^{L_e} n_0 \psi_2(x) dx$$



$$\frac{EA}{L_e} (u_2 - u_1) = N_2 + \int_0^{L_e} n_0 \psi_2(x) dx$$

FE Matrix Equation

$$\textcircled{1} \quad \frac{EA}{L_e}(u_1 - u_2) = -N_1 + \int_0^{L_e} n_0 \psi_1(x) dx$$

$$\textcircled{2} \quad \frac{EA}{L_e}(u_2 - u_1) = N_2 + \int_0^{L_e} n_0 \psi_2(x) dx$$

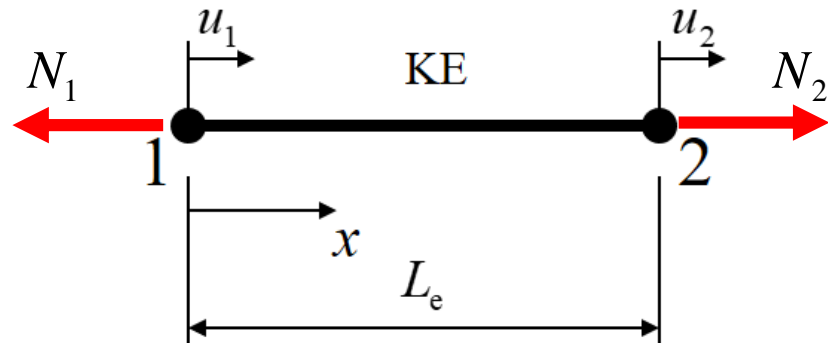


Finite Element Matrix Equation

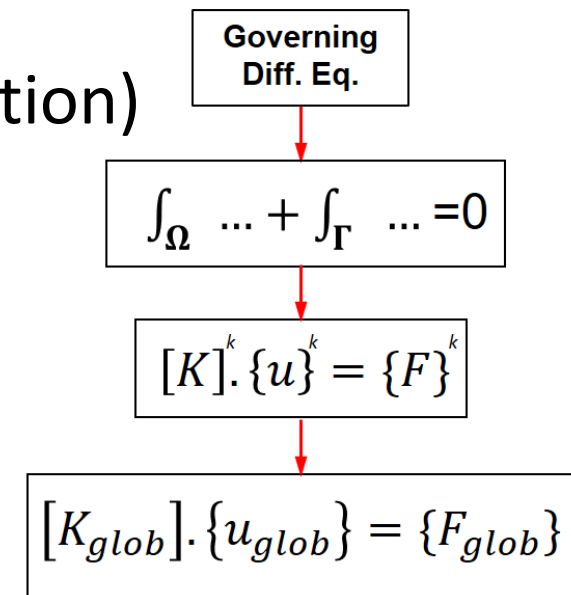
$$\frac{EA}{L_e} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \end{Bmatrix} = \begin{Bmatrix} -N_1 \\ N_2 \end{Bmatrix} + \frac{n_0 L_e}{2} \begin{Bmatrix} 1 \\ 1 \end{Bmatrix}$$

Matrix Notation

$$\mathbf{K}^k \cdot \mathbf{u}^k = \mathbf{f}^k + \mathbf{f}_{n_0}^k$$



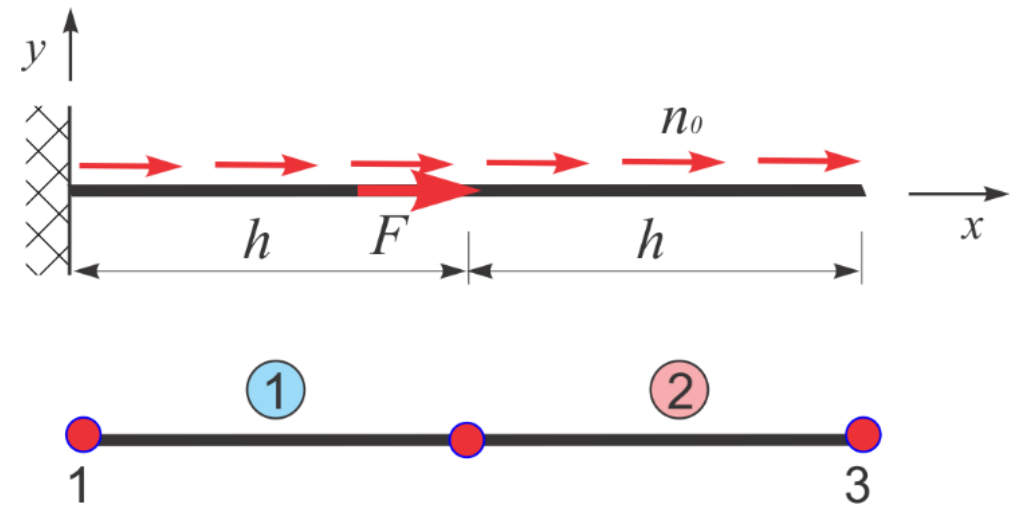
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Meshing & Writing FE Matrix Equation for each FE

$$\frac{EA}{h} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \end{Bmatrix} = \begin{Bmatrix} -N_1^{(1)} \\ N_2^{(1)} \end{Bmatrix} - \frac{n_0 h}{2} \begin{Bmatrix} 1 \\ 1 \end{Bmatrix}$$

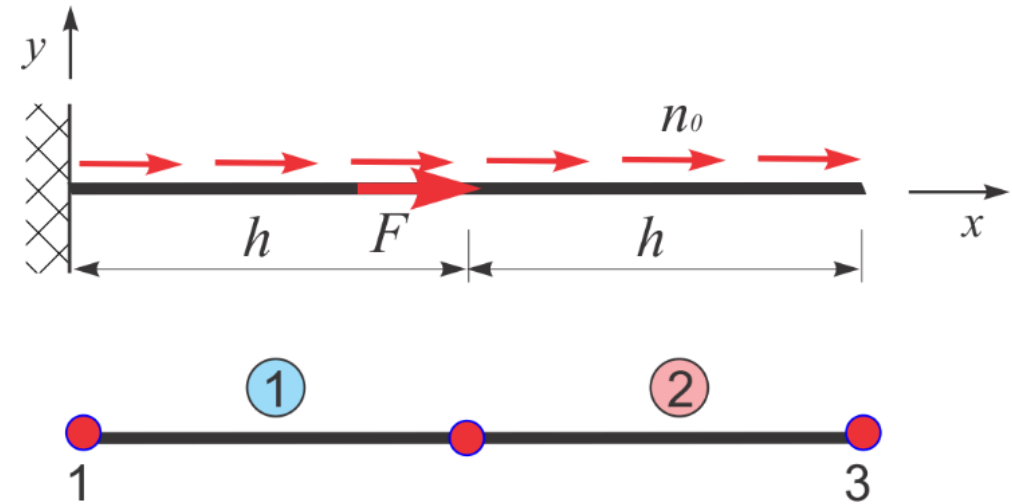
$$\frac{EA}{h} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} u_2 \\ u_3 \end{Bmatrix} = \begin{Bmatrix} -N_2^{(2)} \\ N_3^{(2)} \end{Bmatrix} - \frac{n_0 h}{2} \begin{Bmatrix} 1 \\ 1 \end{Bmatrix}$$



Expansion to all degrees of freedom

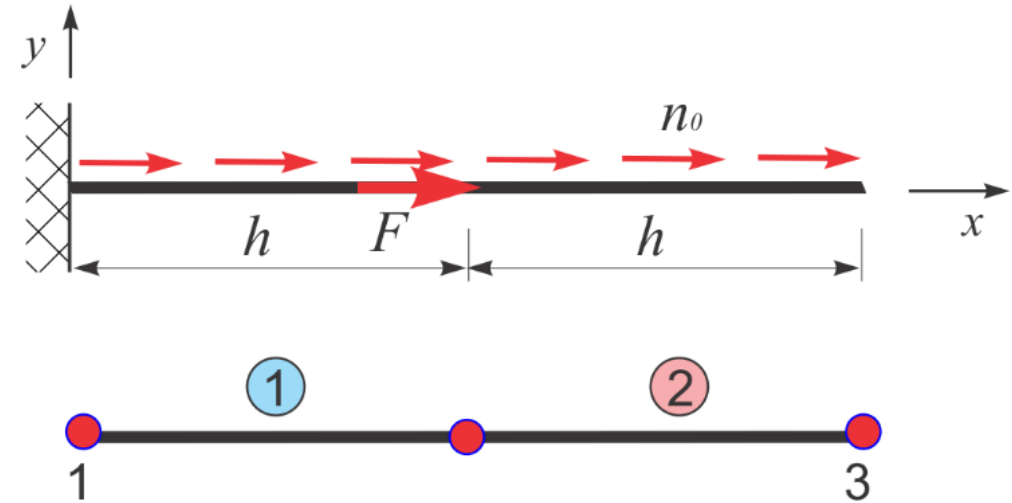
$$\frac{EA}{h} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \end{Bmatrix} = \begin{Bmatrix} -N_1^{(1)} \\ N_2^{(1)} \\ 0 \end{Bmatrix} - \frac{n_0 h}{2} \begin{Bmatrix} 1 \\ 1 \\ 0 \end{Bmatrix}$$

$$\frac{EA}{h} \begin{bmatrix} 0 & 0 & 0 \\ 0 & 1 & -1 \\ 0 & -1 & 1 \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \end{Bmatrix} = \begin{Bmatrix} 0 \\ -N_2^{(2)} \\ N_3^{(2)} \end{Bmatrix} - \frac{n_0 h}{2} \begin{Bmatrix} 0 \\ 1 \\ 1 \end{Bmatrix}$$



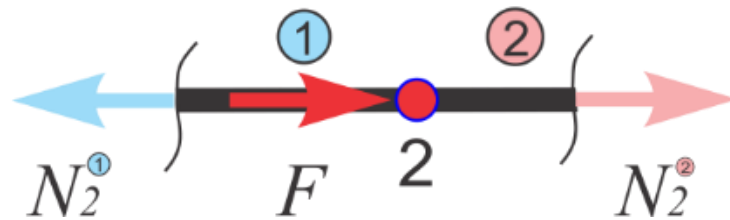
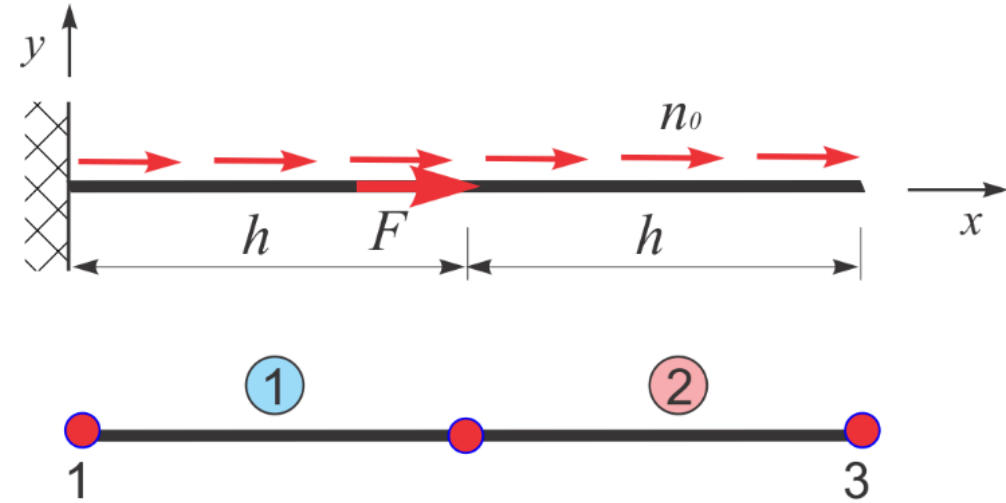
Adding equations together

$$\frac{EA}{h} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 1+1 & -1 \\ 0 & -1 & 1 \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \end{Bmatrix} = \begin{Bmatrix} -N_1^{\textcircled{1}} \\ N_2^{\textcircled{1}} - N_2^{\textcircled{2}} \\ N_3^{\textcircled{2}} \end{Bmatrix} - \frac{n_0 h}{2} \begin{Bmatrix} 1 \\ 1+1 \\ 1 \end{Bmatrix}$$



Adding equations together

$$\frac{EA}{h} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 1+1 & -1 \\ 0 & -1 & 1 \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \end{Bmatrix} = \begin{Bmatrix} -N_1^{(1)} \\ N_2^{(1)} - N_2^{(2)} \\ N_3^{(2)} \end{Bmatrix} - \frac{F_0 n_0 h}{2} \begin{Bmatrix} 1 \\ 1+1 \\ 1 \end{Bmatrix}$$

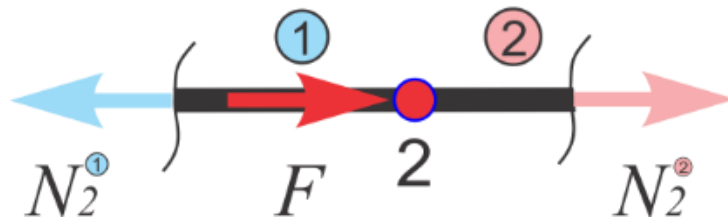
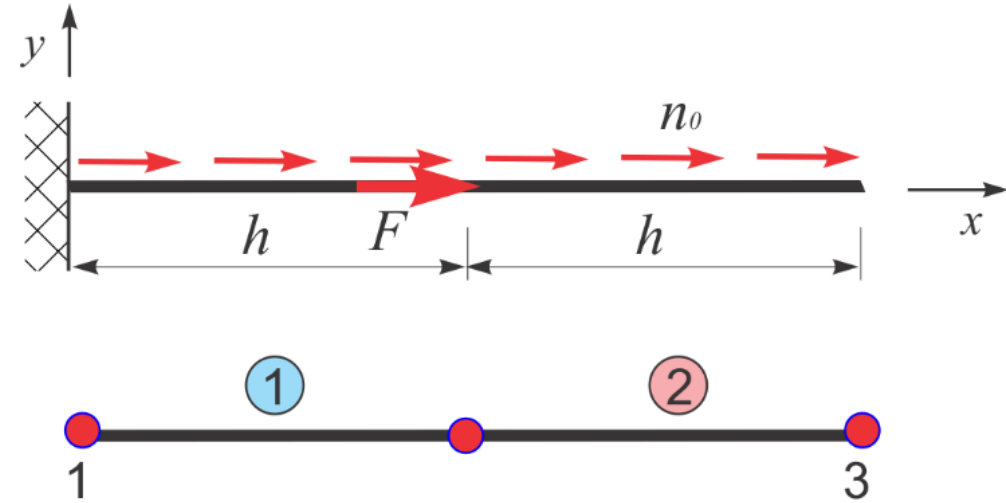


FEM – Boundary conditions

Applying boundary conditions

$$\frac{EA}{h} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 1+1 & -1 \\ 0 & -1 & 1 \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \end{Bmatrix} = \begin{Bmatrix} -N_1^{(1)} \\ N_2^{(1)} - N_2^{(2)} \\ N_3^{(2)} \end{Bmatrix} - \frac{F_0 n_0 h}{2} \begin{Bmatrix} 1 \\ 1+1 \\ 1 \end{Bmatrix}$$

Note: In the original image, the first row of the matrix is highlighted in blue, the second and third rows are highlighted in orange, and the first and third entries of the load vector are marked with a red '0' and crossed out with a red arrow.

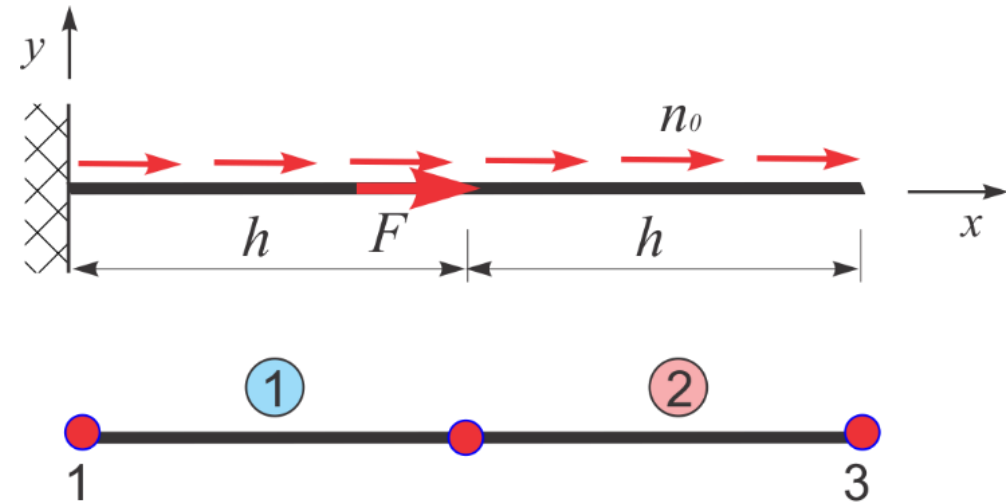


System of equations

$$\frac{EA}{h} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 1 \end{bmatrix} \begin{Bmatrix} 0 \\ u_2 \\ u_3 \end{Bmatrix} = \begin{Bmatrix} -N_1^{\textcircled{1}} \\ F_0 \\ 0 \end{Bmatrix} - \frac{n_0 h}{2} \begin{Bmatrix} 1 \\ 2 \\ 1 \end{Bmatrix}$$

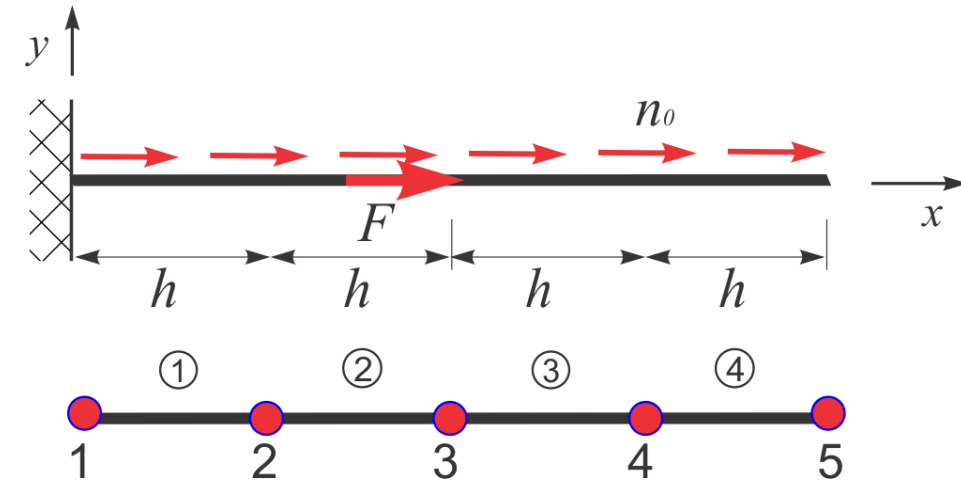
Matrix Notation

$$\mathbf{K}^{glob} \cdot \mathbf{u}^{glob} = \mathbf{f}^{glob} + \mathbf{f}_{n_0}^{glob}$$

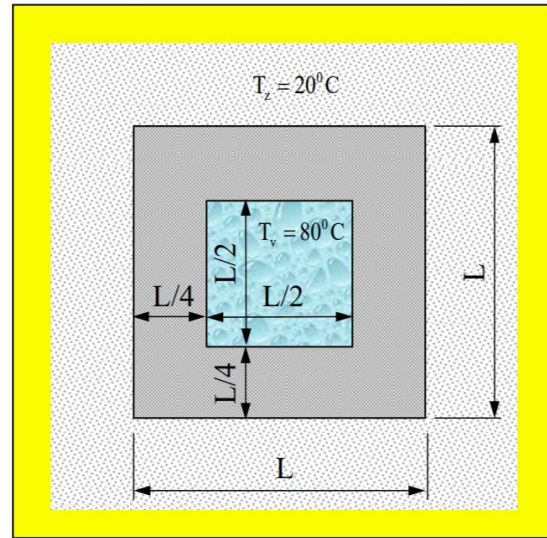


System of equations

$$\frac{EA}{h} \begin{bmatrix} 1 & -1 & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 \\ 0 & -1 & 2 & -1 & 0 \\ 0 & 0 & -1 & 2 & -1 \\ 0 & 0 & 0 & -1 & 1 \end{bmatrix} \begin{Bmatrix} 0 \\ u_2 \\ u_3 \\ u_4 \\ u_5 \end{Bmatrix} = \begin{Bmatrix} -N_1^{(1)} \\ 0 \\ F \\ 0 \\ 0 \end{Bmatrix} - \frac{n_0 h}{2} \begin{Bmatrix} 1 \\ 2 \\ 2 \\ 2 \\ 1 \end{Bmatrix}$$

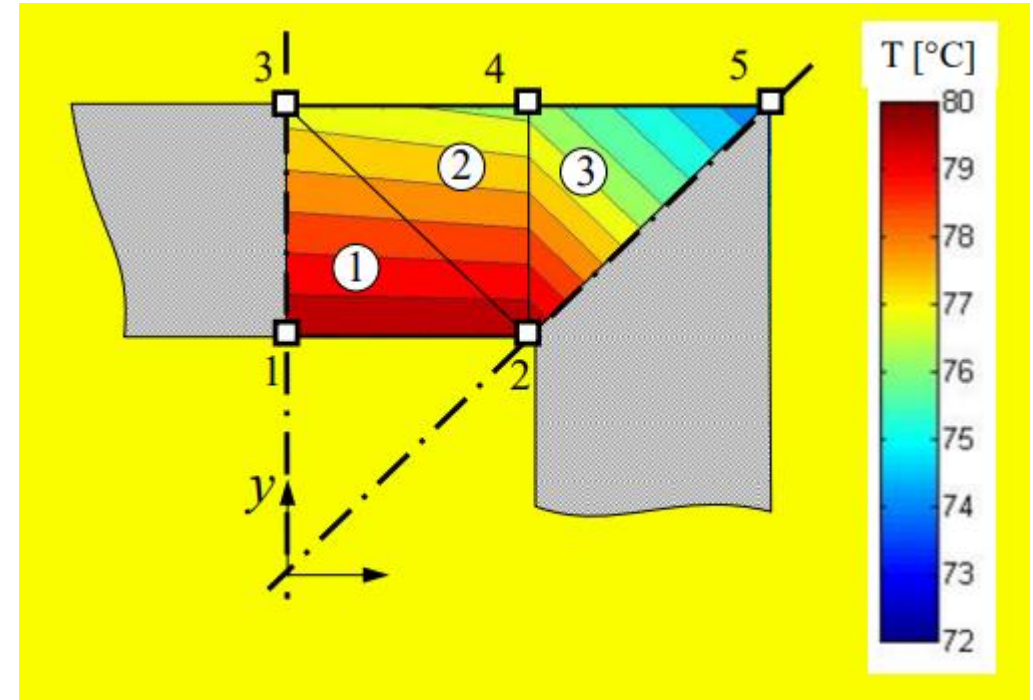
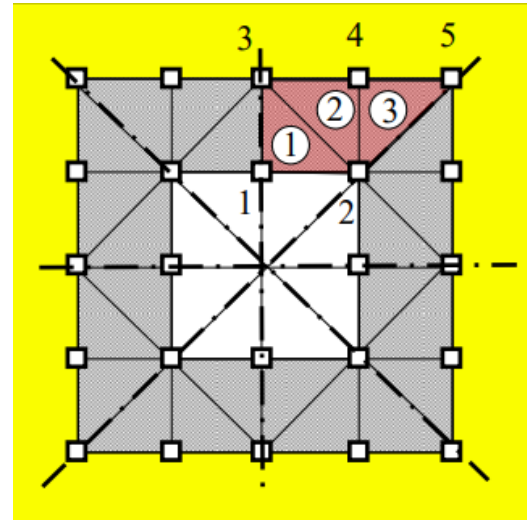


- Stationary heat transfer in 2D



MNM: XIV/19/3f

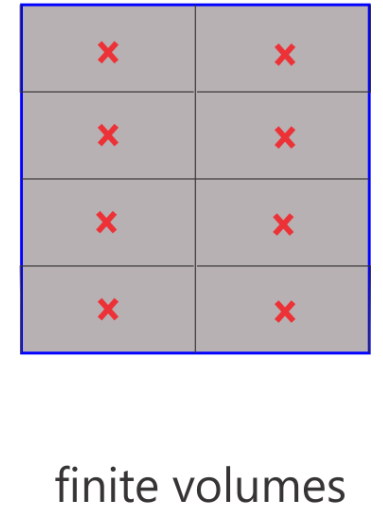
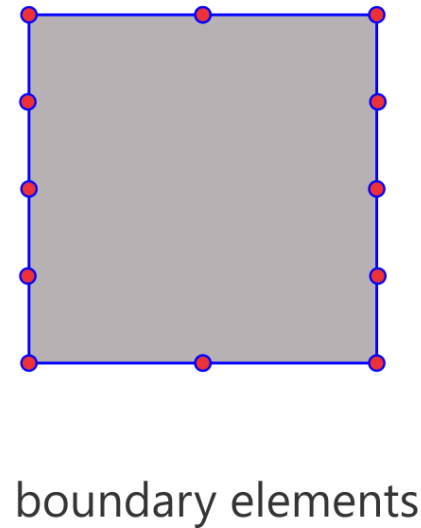
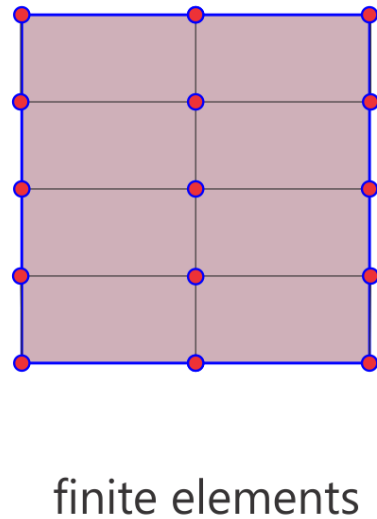
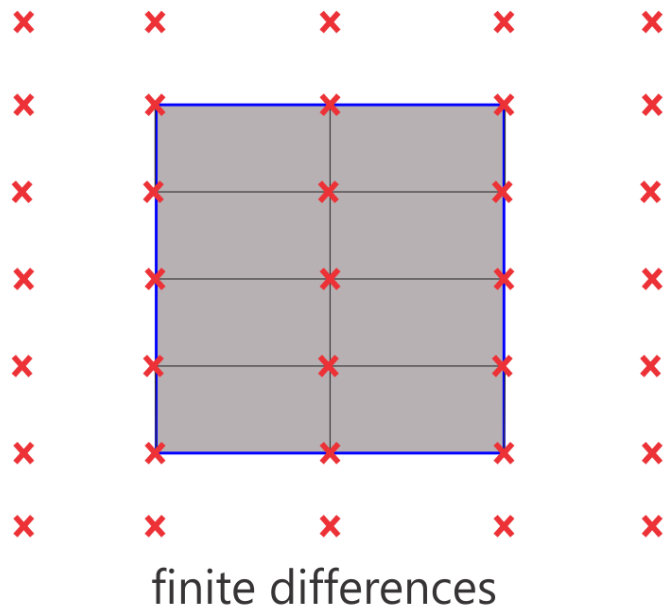
$$k \Delta T(x, y) + q_v = 0$$



- IDEA: integral based approach
- Starting point of derivation is Inverse Integral Form of governing equation
- Solution boundary is discretized into sub-domains, called boundary elements (BEs)
- In the BE sub-domain (only boundary!) we approximate unknown quantities

- ADVANTAGE: solution of governing field equations is converted into searching unknown quantities on boundary
- Suitable for solution of potential problems & infinite problems
- DISADVANTAGE: full and non-symmetrical matrices

- Basic differences between BVP solution methods



- IDEA: integral based approach
- Starting point of derivation is Integral Form of governing equation, where the integral over solution domain is converted in surface integral around the domain
- Solution domain is discretized into sub-domains, called finite volumes (cells or control volumes)
- The unknown solution variable is constant over the cell (calculated at the centre)

- ADVANTAGE: solution of governing field equations is converted into searching unknown quantity in the cell centre.
- Continuity requirements are simply fulfilled.
- Similar to FDM, Suitable for solution of heat transfer & fluid flow problems
- DISADVANTAGE: fulfilment of boundary conditions related to solution variable

- Presentation on the stationary heat transfer problem

$$\frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left(k \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left(k \frac{\partial T}{\partial z} \right) + q_v = 0$$

$$\text{div} [k \text{ grad}(T)] = -q_v$$

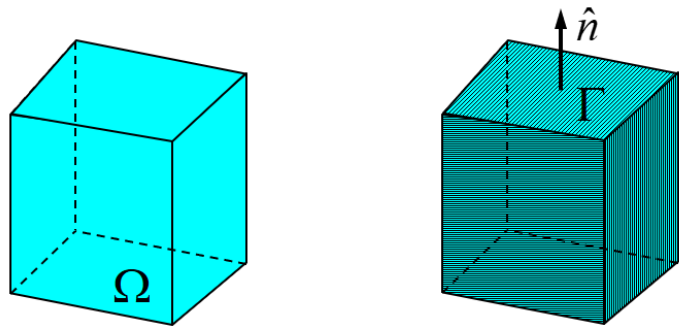


$$\int_{\Omega} \text{div} [k \text{ grad}(T)] d\Omega = - \int_{\Omega} q_v d\Omega$$

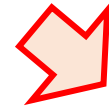
Integration over the domain

- Green-Gauss theorem

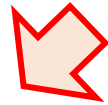
$$\int_{\Omega} \text{div}[k \text{ grad}(T)] d\Omega = \int_{\Gamma} k \text{ grad}(T) \hat{n} d\Gamma$$



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$$\int_{\Gamma} k \text{ grad}(T) \hat{n} d\Gamma = -\int_{\Omega} q_V d\Omega$$

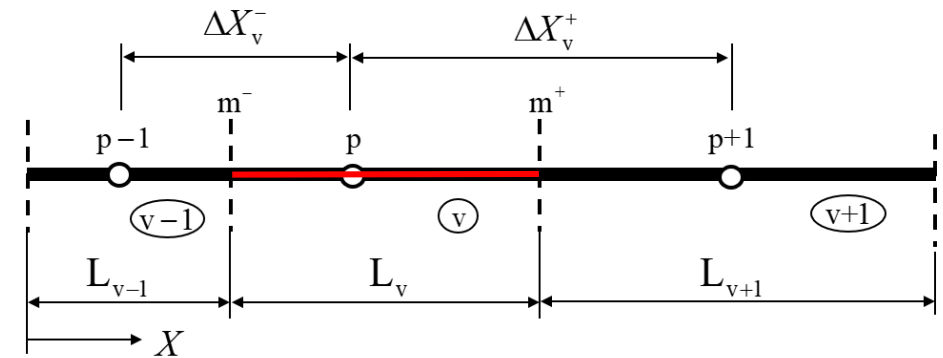
Heat balance over control volume

- Reduction to 1D $T = T(x)$

$$\int_{\Gamma} k \operatorname{grad}(T) \hat{n} d\Gamma = - \int_{\Omega} q_v d\Omega \quad \Rightarrow \quad \int_{\Gamma} k \frac{dT}{dx} n_x d\Gamma = - \int_0^L q_v A dx$$

Since the temperature is unknown, we use approximation

$$\int_{\Gamma_v} k \frac{dT}{dx} n_x d\Gamma \approx \left[\left(\frac{k A}{\Delta X_v} \right)_{m^+} (T_{p+1} - T_p) \right] - \left[\left(\frac{k A}{\Delta X_v} \right)_{m^-} (T_p - T_{p-1}) \right]$$



- Reduction to 1D $T = T(x)$

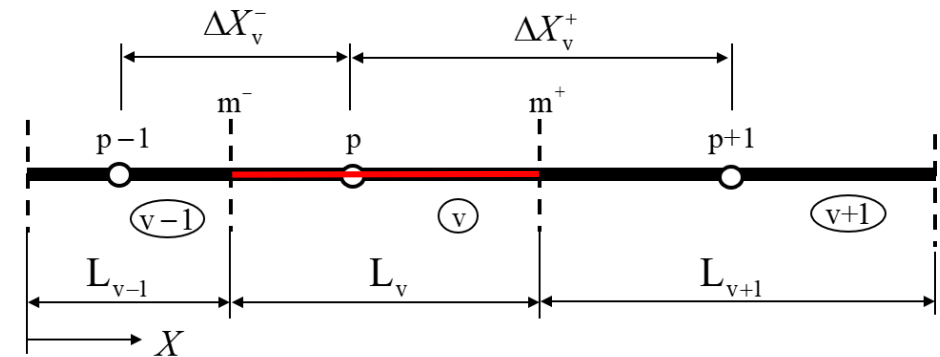
$$\int_{\Gamma} k \operatorname{grad}(T) \hat{n} d\Gamma = - \int_{\Omega} q_V d\Omega \quad \Rightarrow \quad \int_{\Gamma} k \frac{dT}{dx} n_x d\Gamma = - \int_0^L q_V A dx$$

Right-hand side of the equation, we approximate as

$$- \int_0^{L_v} q_V A dx \approx - (q_V)_p A_p L_v$$

Finite volume equation

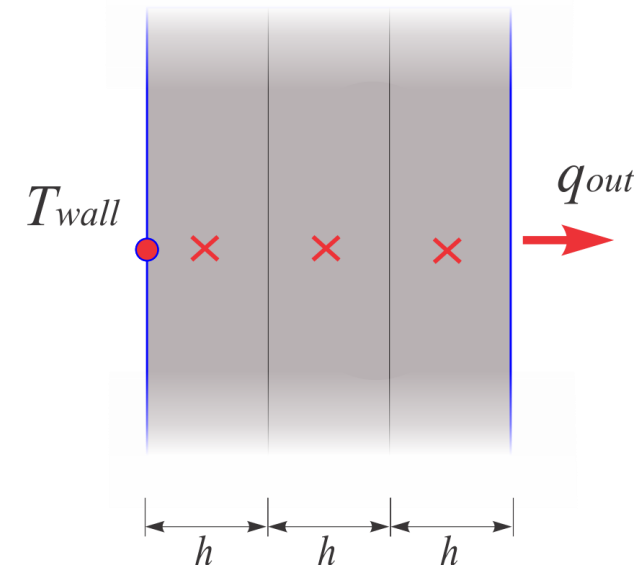
$$\left[\left(\frac{kA}{\Delta X_v} \right)_{m^+} (T_{p+1} - T_p) \right] - \left[\left(\frac{kA}{\Delta X_v} \right)_{m^-} (T_p - T_{p-1}) \right] = - (q_V)_p A_p L_v$$



- System of equations

$$\frac{k}{h} \begin{bmatrix} -3 & 1 & 0 \\ 1 & -2 & 1 \\ 0 & 1 & -1 \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \\ T_3 \end{Bmatrix} = \begin{Bmatrix} -2k/hT_{wall} \\ 0 \\ -q_{out} \end{Bmatrix}$$

Heat transfer problem



Thank you for your attention!

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